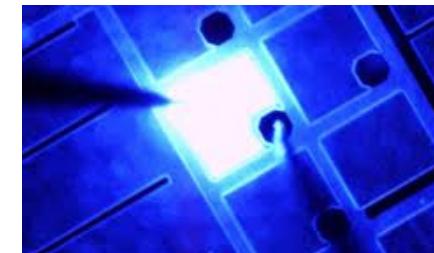


Lecture 9 – 16/04/2025

Light-emitting diodes

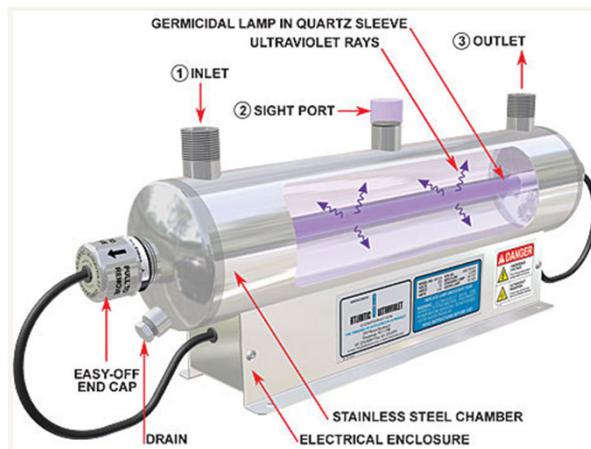
- Basic properties
- Notion of efficiency

Chap. 13 in Rosencher-Vinter ≡
reference chapter until Lecture 14!

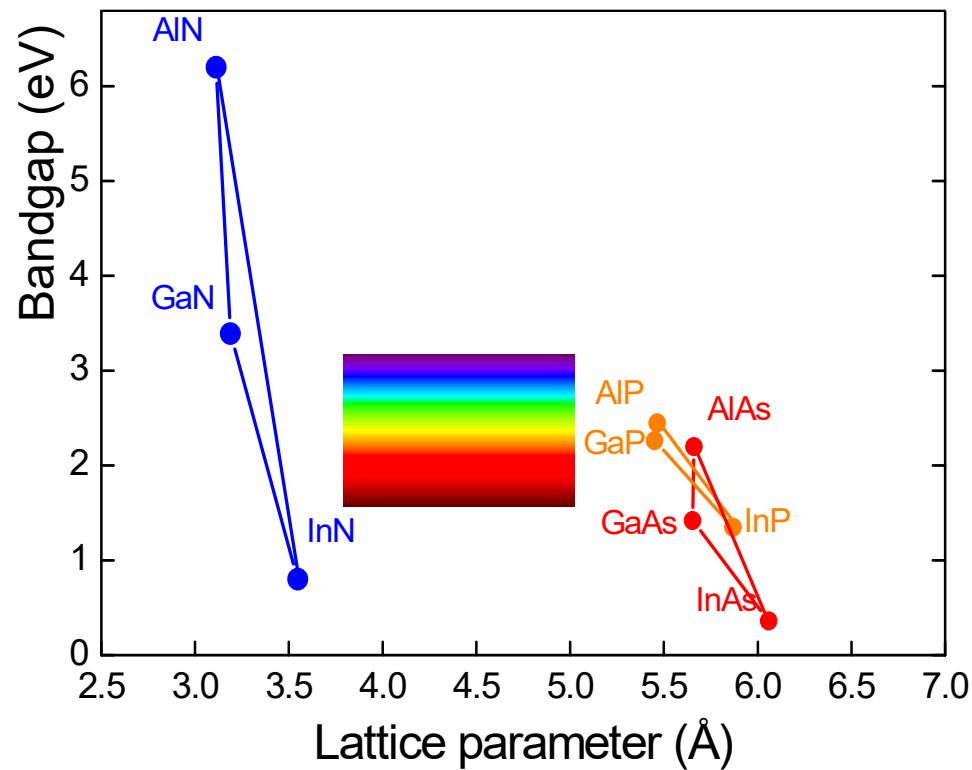


Main applications of light-emitting diodes

- Displays
- Lighting
- Communication
- Purification (UV)
- 3D sensing



Semiconductors for optoelectronics

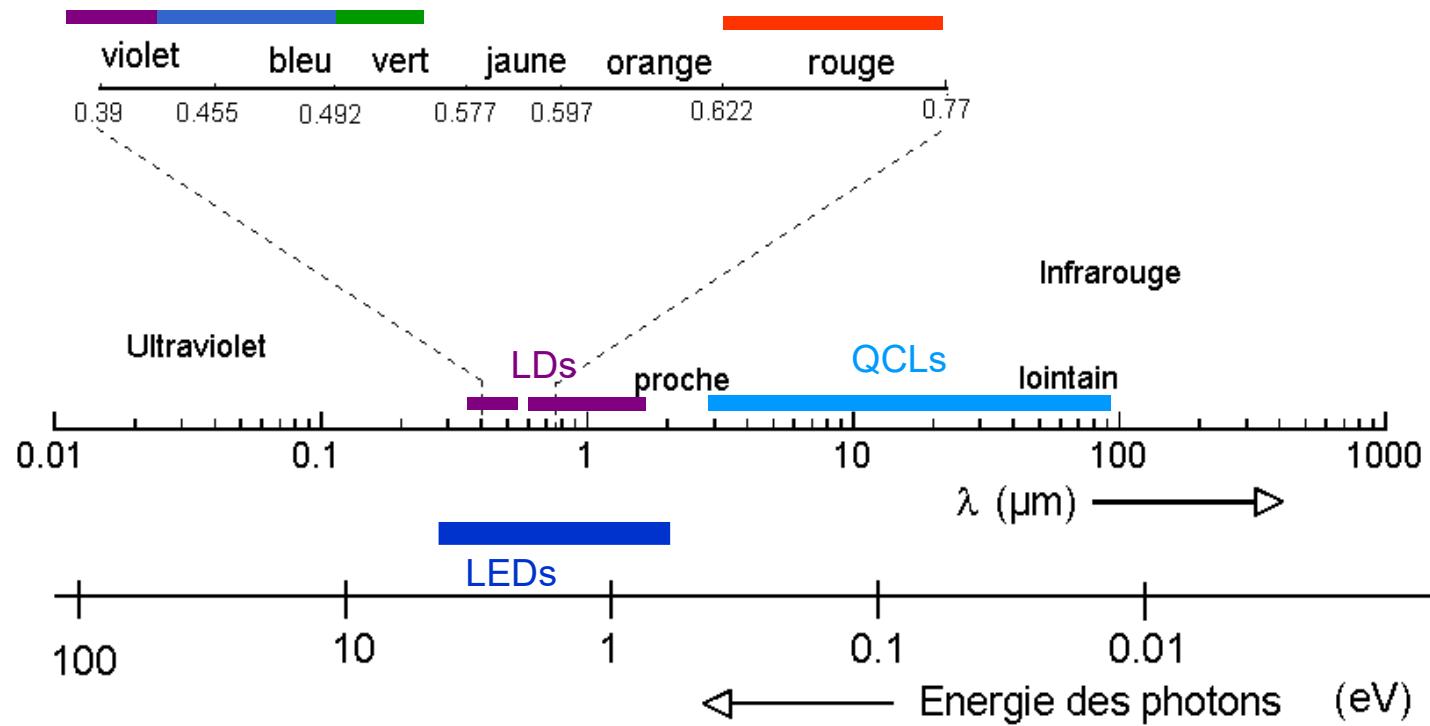


Arsenides: $(Al,Ga,In)As$

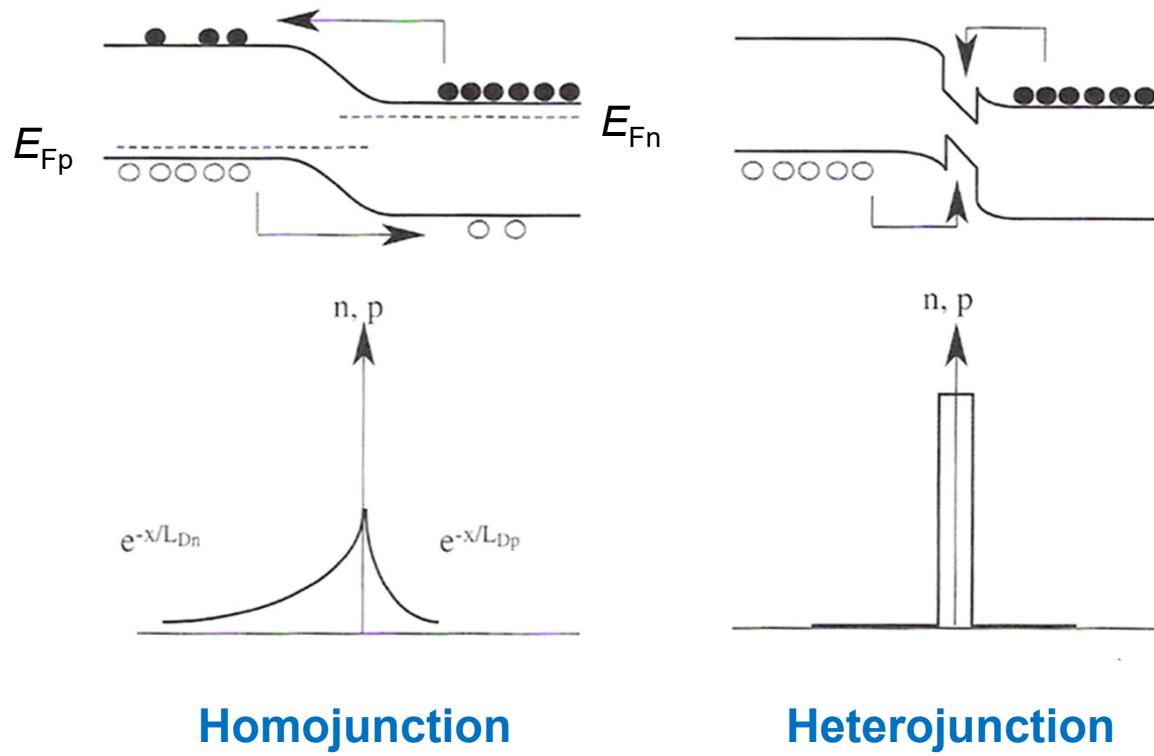
Phosphides: $(Al,Ga,In)P$

Nitrides: $(Al,Ga,In)N$

Spectral domain covered by commercial LEDs and LDs



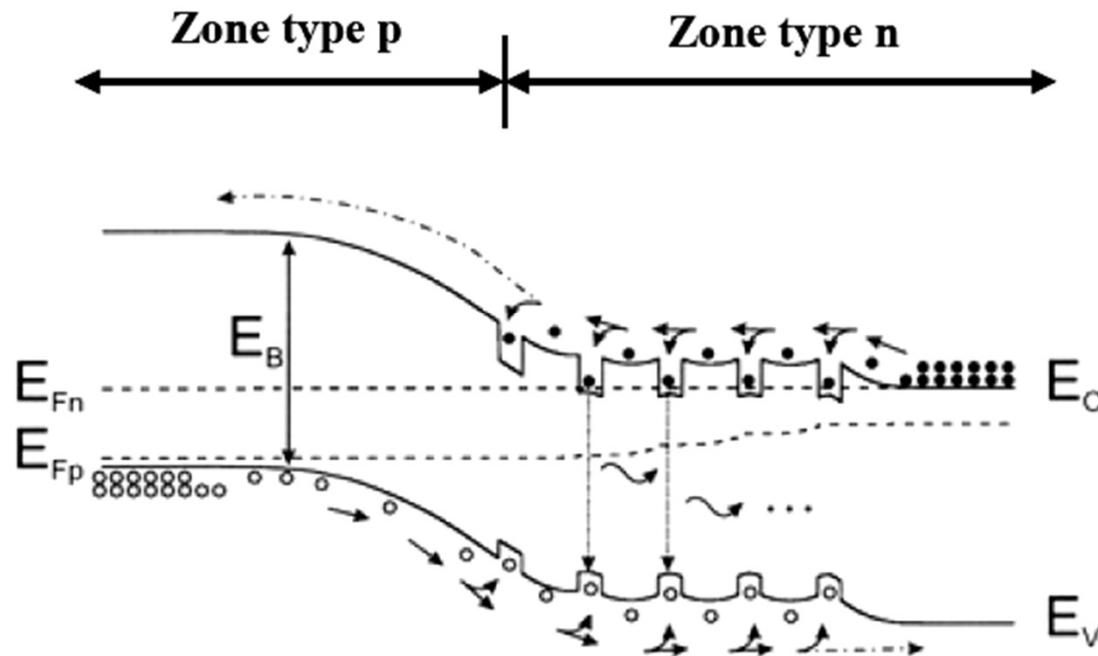
LED structures



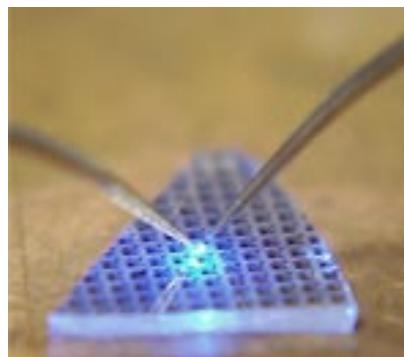
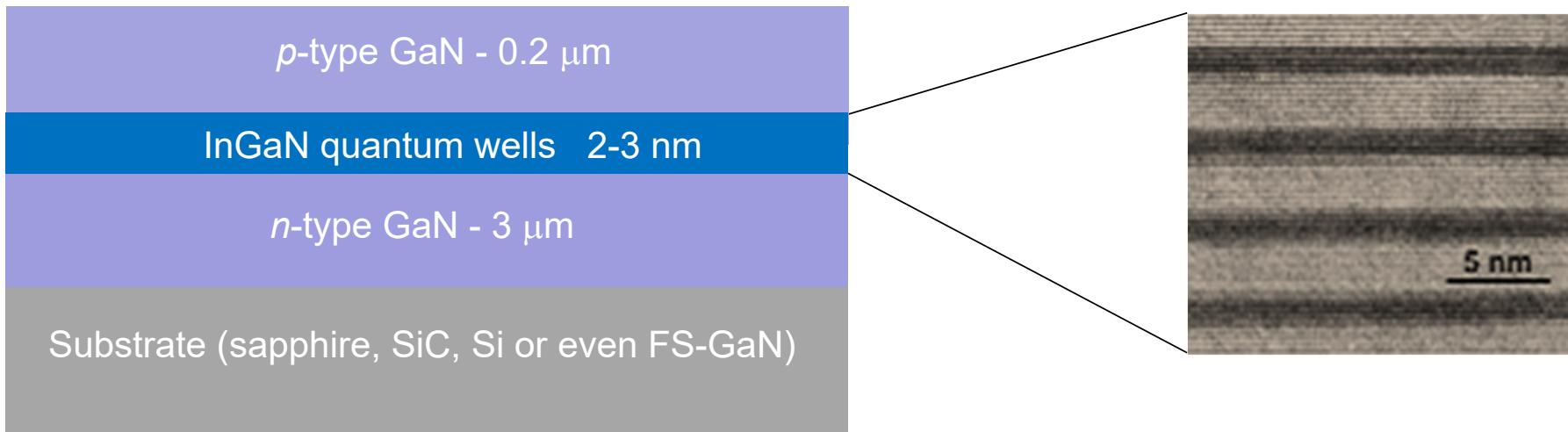
Heterojunction allows for an efficient spatial trapping of injected carriers \Rightarrow increased radiative efficiency, improved operating characteristics (L - I & I - V curves)

LED structures

Multiple quantum well LED



Blue LED structure: a basic picture



Substrate \Rightarrow epilayer material quality

***n*-type and *p*-type doped layers** \Rightarrow efficient injection

Active region \Rightarrow radiative efficiency

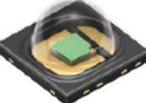
Key markets for III-N LEDs

- **Automotive**
 - Forward lighting
 - 3D sensing
- **Consumer**
 - Projection
 - Tablets/monitors/TV
- **Industry**
 - Video walls
 - White goods
 - 3D sensing
- **General lighting**
 - Indoor/outdoor lighting
 - Shop lighting



Sources: OSRAM Opto Semiconductors

Emerging LED market: toward 3D sensing

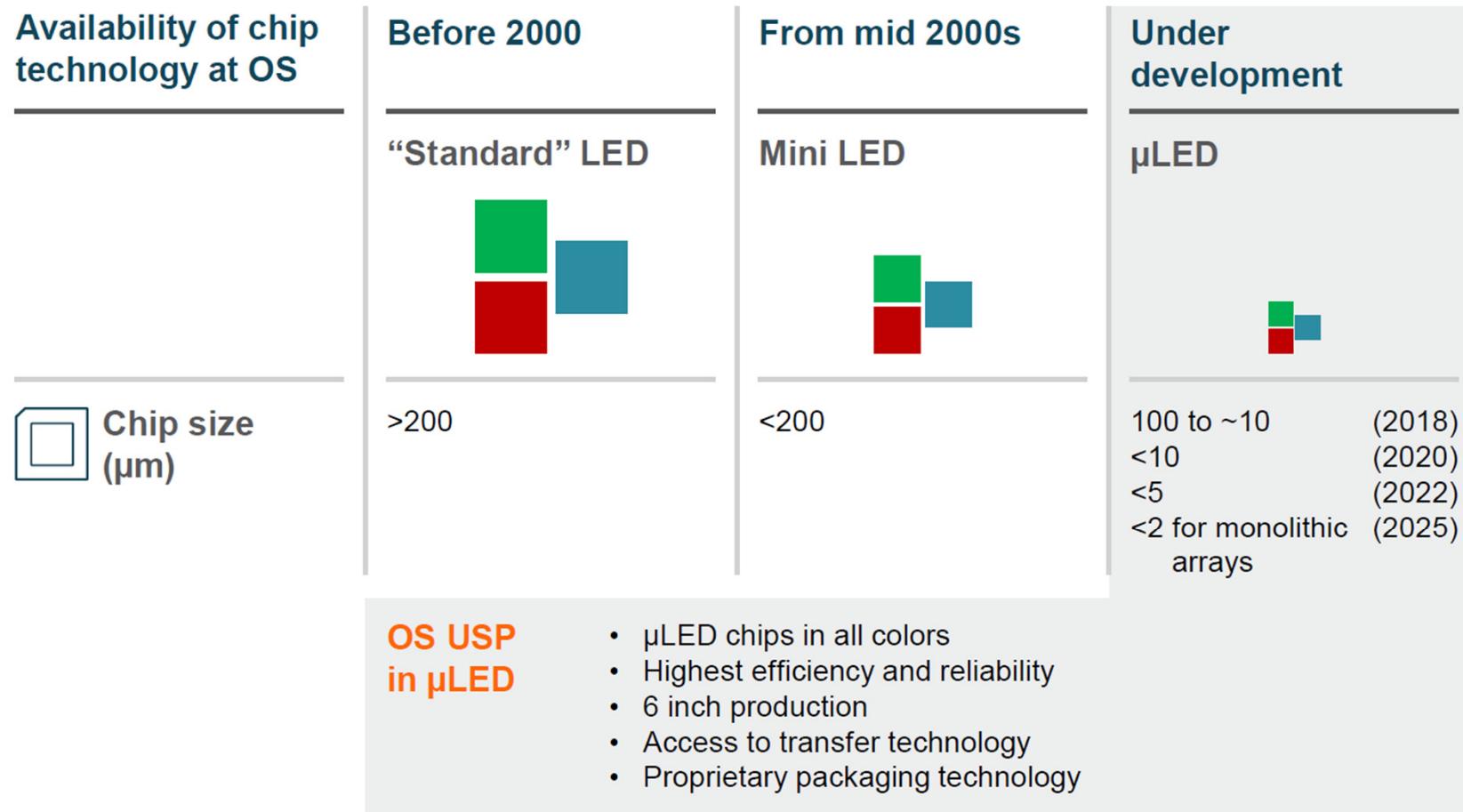
3D Sensing applications		Emitter technologies		
Examples		LED	Edge emitter	VCSEL ²⁾
Mobile Devices	 			
Industry	 			Deep-Dive
Auto-motive ¹⁾	 			

1) Different market that shows same emitter technology

2) Vertical-cavity surface-emitting laser

LED chips: current trends

Evolution of chip sizes

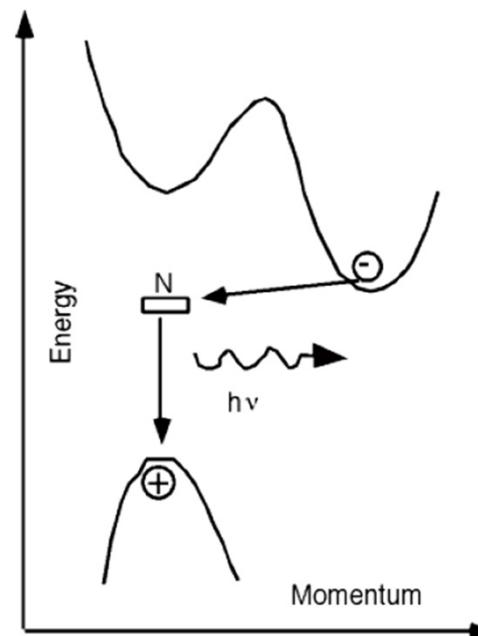
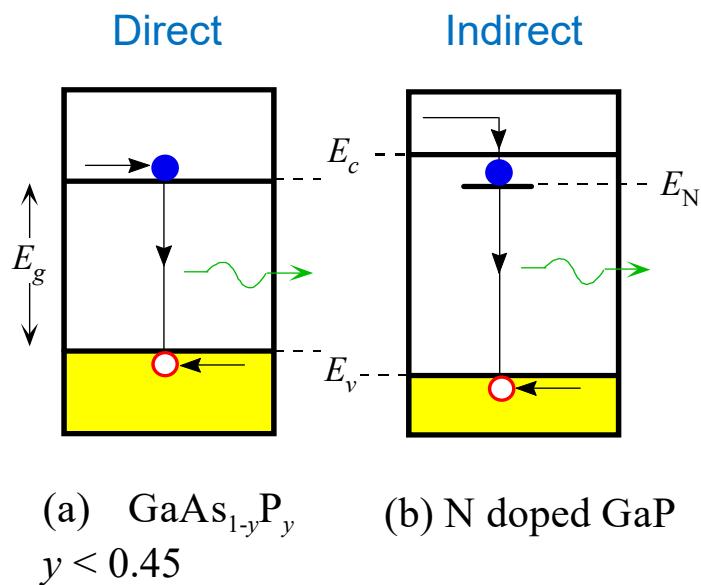


LED market

Projection applications	Emitter technologies				Not used
	LED	High power blue Laser	RGB low power laser	Pixelated μ LED array	
Home Projection 	✓				
Professional Projection 	✓	✓			
Mobile Projection 	✓		✓		
Augmented Reality 	✓			✓	
					Future scope

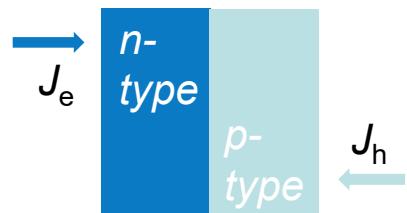
LEDs: basic properties

LEDs made of indirect bandgap semiconductors

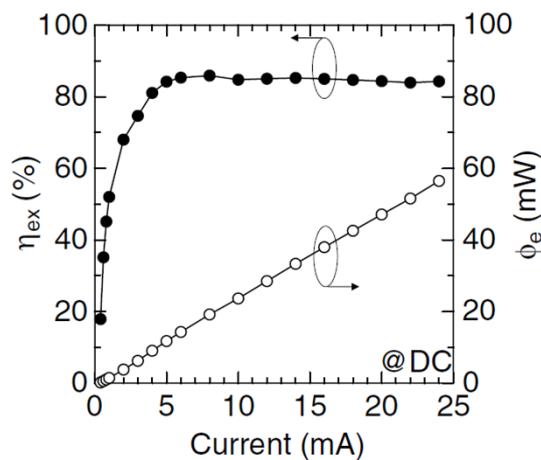


Indirect bandgap \Rightarrow luminescence through a localized defect lying within the bandgap

LEDs: basic properties



Cf. Lectures 7 & 14, fall semester + Chap. 7 Rosencher-Vinter



$$\frac{J_e}{qd} = \frac{J_h}{qd} = \frac{J}{qd} = \frac{n}{\tau_{\text{tot}}} = A_{\text{nr}}n + Bn^2 + C_{\text{Aug}}n^3$$

$d = V/S$ Thickness of the active region

$$\frac{1}{\tau_{\text{nr}}} = A_{\text{nr}} + C_{\text{Aug}}n^2$$

$$\frac{1}{\tau_r} = Bn$$

$$\frac{1}{\tau_{\text{tot}}} = \frac{1}{\tau_{\text{nr}}} + \frac{1}{\tau_r}$$

Stimulated emission term neglected

Out of equilibrium carrier density

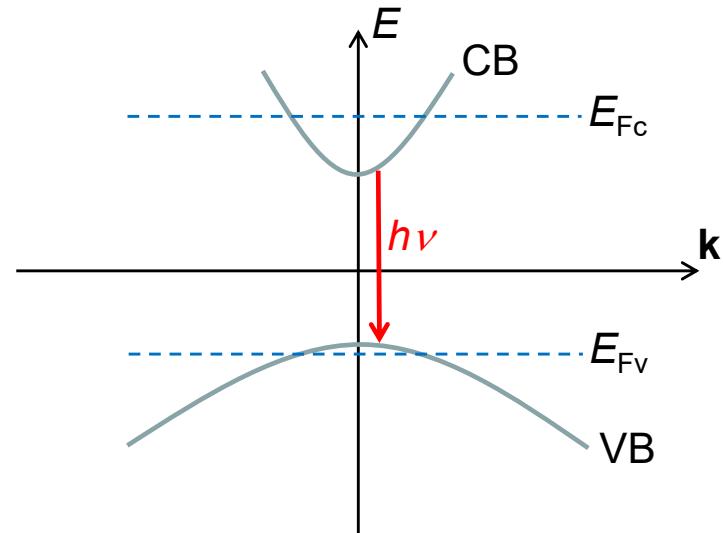
$$n = \frac{J\tau_{\text{tot}}}{qd}$$

⇒ Strong dependence on the thickness of the active region (homo- vs heterojunction (QWs, etc.))

Once n is known, possibility to derive the position of the quasi-Fermi levels E_{F_n} and E_{F_p}

Electrical injection

- Both the valence and the conduction bands get more and more filled upon increasing current injection
- The carrier populations are described by the quasi-Fermi levels E_{Fc} and E_{Fv}



$$f_c(E) = \frac{1}{\exp\left(\frac{E - E_{Fc}}{k_B T}\right) + 1}$$

$$f_v(E) = \frac{1}{\exp\left(\frac{E - E_{Fv}}{k_B T}\right) + 1}$$

Note that here $f_v(E)$ describes the evolution of the electron population in the valence band!

Electrical injection

Determination of the quasi-Fermi level

$$n = \int_{E_c}^{\infty} \frac{1}{\exp\left(\frac{E - E_{F_c}}{k_B T}\right) + 1} \rho_c(E) dE$$

How many photons are emitted?

$$\xrightarrow{J_e} \frac{n_e}{\text{CB}}$$

$$J_e = J_h = J \quad \text{electrical neutrality}$$

and $n_e = n_h = n$ (if the doping levels are not too high)

Steady-state \Rightarrow recombination in the active region

$$\frac{n_h}{\text{VB}} \quad J_h$$

←

The number of emitted photons is then given by

$$R_{\text{tot}} \times \text{Volume} = J/q \times S \quad \text{← Contact size}$$

with R_{tot} the recombination rate (per unit volume)

Electrical injection

Different paths for electron-hole recombinations

- Non-radiative

$$A_{nr} n$$

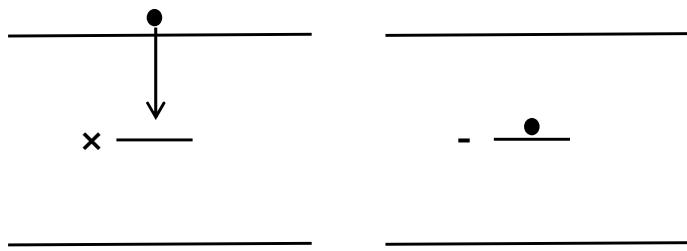
- Spontaneous

$$B n^2$$

B bimol. coeff. $\sim 10^{-12}\text{-}10^{-10} \text{ cm}^3\text{s}^{-1}$

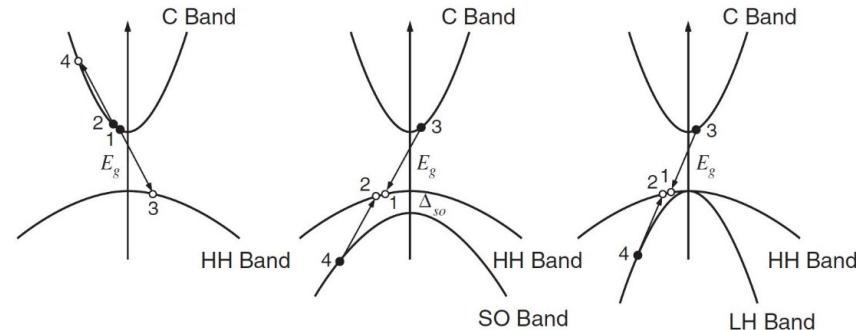
- Auger

$$C n^3$$



Shockley-Read-Hall
recombinations

Cf. Lecture 7, fall semester + Chaps. 5-7
Rosencher-Vinter



Auger
recombinations

Quasi-Fermi levels in bulk semiconductors (reminder)

Example: quasi-Fermi levels in bulk GaAs

Non-degenerate case

$$E_{F_n} = E_C - k_B T \ln \left(\frac{N_C}{n} \right)$$

$$E_{F_p} = E_V + k_B T \ln \left(\frac{N_V}{p} \right)$$

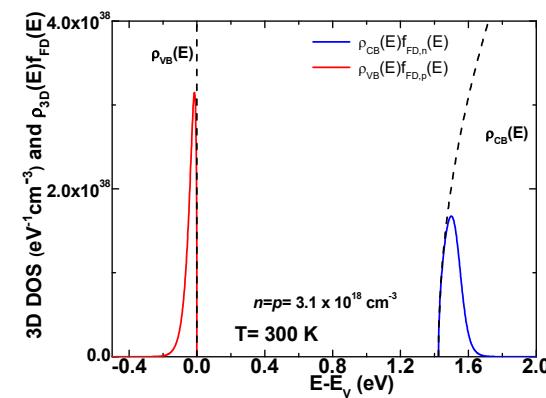
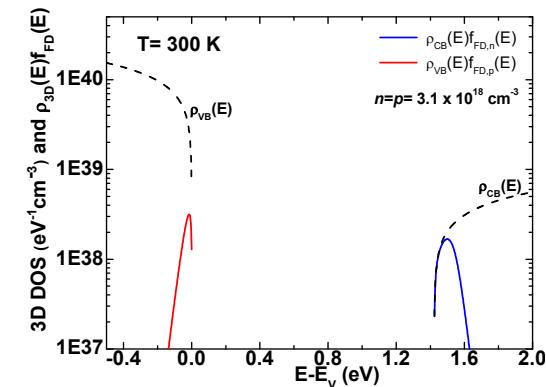
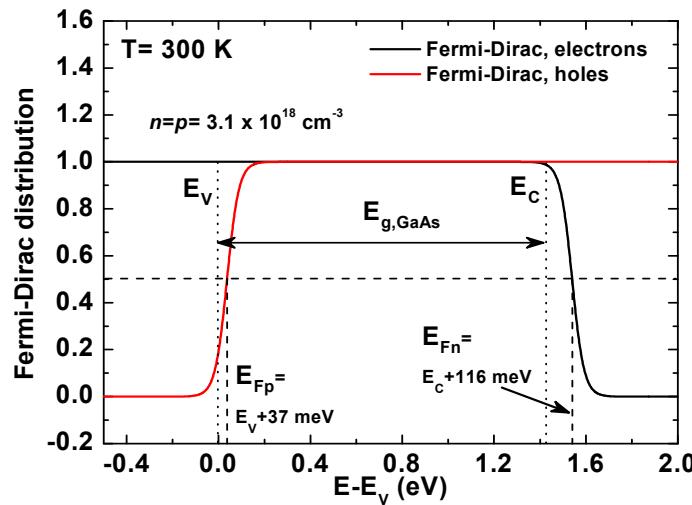
Degenerate case

$$E_{F_n} = E_C + \frac{\hbar^2}{2m_C^*} \left(3\pi^2 n \right)^{2/3}$$

$$E_{F_p} = E_V - \frac{\hbar^2}{2m_V^*} \left(3\pi^2 p \right)^{2/3}$$

effective DOS

$$N_{C,V} = \frac{1}{4} \left(\frac{2m_{C,V}^* k_B T}{\pi \hbar^2} \right)^{3/2}$$



Spontaneous emission

In an intrinsic bulk semiconductor:

Cf. Lectures 12 & 14, fall semester
+ Chap. 7 Rosencher-Vinter

The spontaneous recombination rate (s^{-1}) between the CB and the VB is given for a state with a wavevector \mathbf{k}

$$r_{\text{sp}}(\mathbf{k}) = A_{\text{cv}} f_{\text{c}}(E_{\text{c}}(\mathbf{k}))(1-f_{\text{v}}(E_{\text{v}}(\mathbf{k})))$$

with $A_{\text{CV}} = 1/\tau_{\text{R}}$ the spontaneous recombination rate

and the radiative lifetime is given by

$$\tau_{\text{R}} = \frac{\pi c^3 \hbar \epsilon_0}{q^2 x_{\text{vc}}^2 n_{\text{op}} \omega_{\text{vc}}^3} = \frac{2\pi c^3 \hbar^2 \epsilon_0 m_0}{q^2 n_{\text{op}} E_{\text{g}} E_{\text{P}}}$$

$\tau_{\text{R}} \uparrow \text{when } E_{\text{g}} \downarrow$

with x_{vc} the interband dipolar optical matrix element and E_{P} the Kane energy ($\sim 20\text{-}22 \text{ eV}$)

⇒ *It is more challenging to achieve a laser based on a wide bandgap SC!*

Spontaneous emission

In an intrinsic bulk semiconductor: Cf. Lecture 14, fall semester + Chap. 7 Rosencher-Vinter

The spectral distribution of spontaneous recombination rate $R_{\text{sp}}(h\nu)$ due to a quasi-equilibrium distribution of carriers is then given by

Spin-related

$$R_{\text{sp}}(h\nu) = 2 \sum_{\mathbf{k}} r_{\text{sp}}(\mathbf{k}) = 2 \sum_{\mathbf{k}} \frac{1}{\tau_{\text{R}}(\mathbf{k})} f_{\text{c}}(\mathbf{k})(1 - f_{\text{v}}(\mathbf{k})) \delta(E_{\text{c}} - E_{\text{v}} = h\nu)$$

The summation is performed over all \mathbf{k} -vectors verifying the energy conservation condition (hence the Dirac delta)

$$E_{\text{c}}(\mathbf{k}) - E_{\text{v}}(\mathbf{k}) = h\nu = E_{\text{g}} + \frac{\hbar^2 k^2}{2m_{\text{r}}} \quad \text{Expression relying on the verticality of optical transitions in } \mathbf{k}\text{-space}$$

which leads to

$$R_{\text{sp}}(h\nu) = \int_0^{\infty} r_{\text{sp}}(E) \rho_{\text{j}}(E) \delta(E = h\nu) dE = r_{\text{sp}}(h\nu) \rho_{\text{j}}(h\nu) \quad \frac{1}{m_{\text{r}}} = \frac{1}{m_{\text{c}}^*} + \frac{1}{m_{\text{v}}^*}$$

$$R_{\text{sp}}(h\nu) = \frac{1}{\tau_{\text{R}}} \rho_{\text{j}}(h\nu) f_{\text{c}}(E_{\text{c}}(h\nu))(1 - f_{\text{v}}(E_{\text{v}}(h\nu)))$$

Joint density of states (JDOS)

LEDs: basic properties

Spectral distribution of spontaneous recombination rate

$$R_{sp}(h\nu) = \frac{1}{\tau_r} \rho_j(h\nu) f_c(h\nu) (1 - f_v(h\nu)),$$

$$f_c(h\nu) = \frac{1}{1 + \exp\left(\frac{E_c(h\nu) - E_{F_n}}{k_B T}\right)} \approx \exp\left(-\frac{E_c(h\nu) - E_{F_n}}{k_B T}\right),$$

Usually valid for LEDs (non-degenerate case, i.e., Boltzmann approximation is valid)

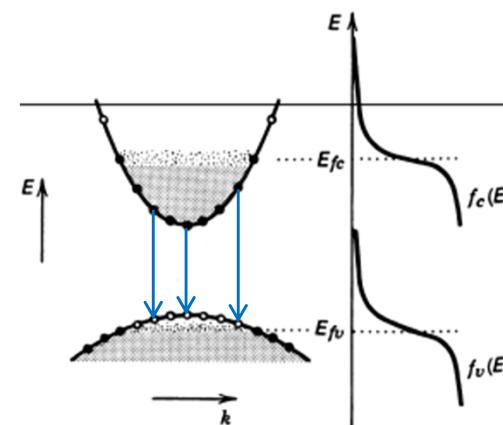
$$E_c(h\nu) = E_g + \frac{m_r}{m_c} (h\nu - E_g),$$

$$f_v(h\nu) = \frac{1}{1 + \exp\left(\frac{E_v(h\nu) - E_{F_p}}{k_B T}\right)},$$

$$E_v(h\nu) = -\frac{m_r}{m_v} (h\nu - E_g),$$

$$\rho_j(h\nu) = \frac{1}{2\pi^2} \left(\frac{2m_r}{\hbar^2}\right)^{3/2} (h\nu - E_g)^{1/2}.$$

Joint density of states (JDOS)

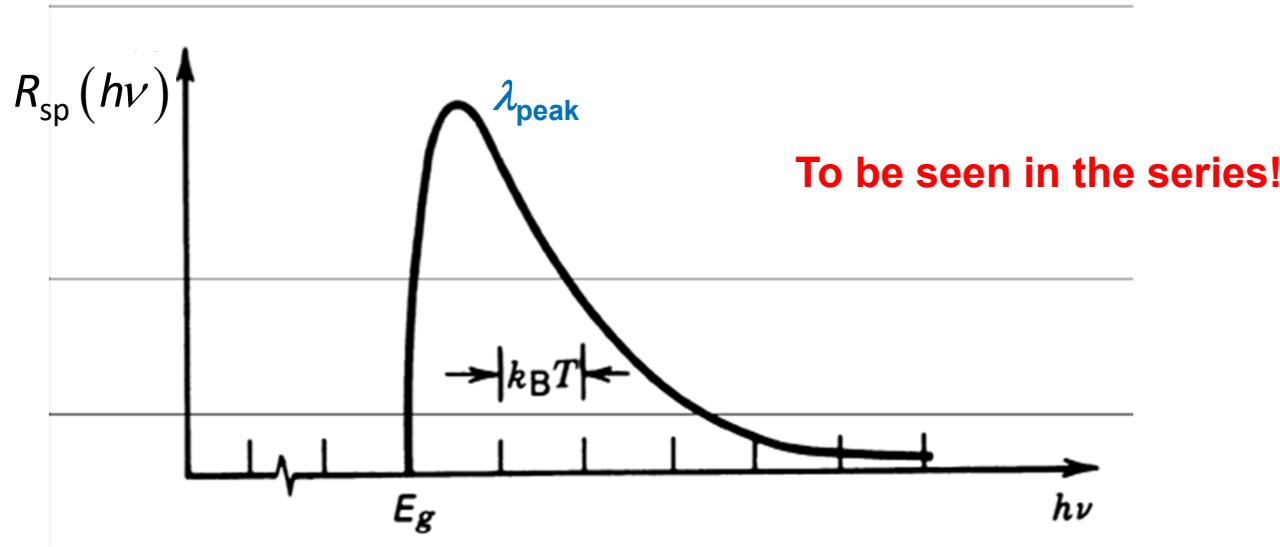


LEDs: basic properties

$$R_{\text{sp}}(h\nu) = K_{\text{sp}} (h\nu - E_g)^{1/2} \exp\left(-\frac{h\nu - E_g}{k_B T}\right),$$

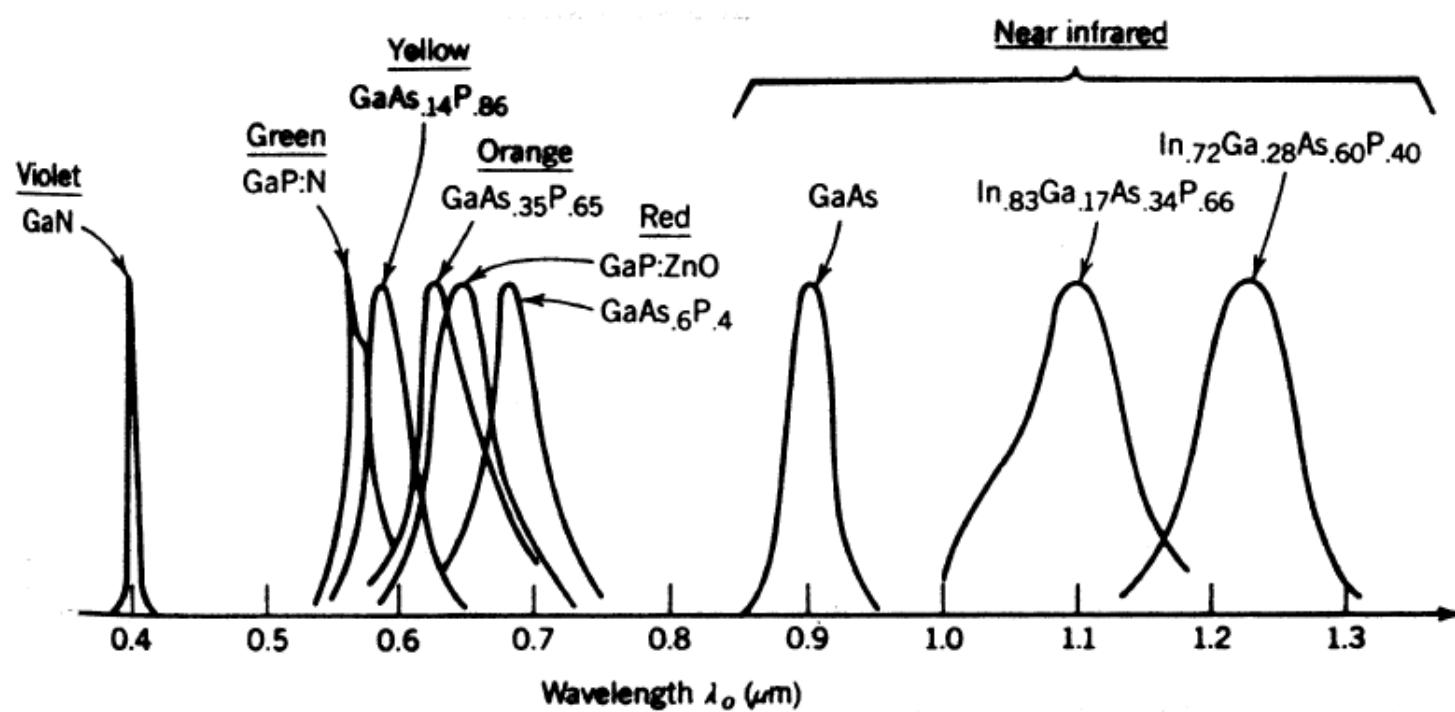
with $K_{\text{sp}} = \frac{(2m_r)^{3/2}}{\pi\hbar^2\tau_r} \exp\left(\frac{\Delta E_F - E_g}{k_B T}\right) = \frac{(2m_r)^{3/2}}{\pi\hbar^2\tau_r} e^{\frac{eV_{\text{app}} - E_g}{k_B T}}$.

$E_{F_C} - E_{F_V}$



LEDs emission spectra

Spectral density emitted by LEDs



$$\Delta\lambda \propto \lambda_{\text{peak}}^2 k_B T$$

LED efficiency

Internal quantum efficiency (IQE, η_i):

$\eta_i = [\text{generated photons}]/[\text{injected electrons}]$ with [] \equiv particle number

$$\eta_i = \frac{\tau_{\text{tot}}}{\tau_r} = \frac{\tau_{\text{nr}}}{\tau_{\text{nr}} + \tau_r} = \frac{Bn}{A_{\text{nr}} + Bn + Cn^2}$$

with τ_{nr} non-radiative lifetime
 τ_r radiative lifetime

The internal quantum efficiency can be as high as 99% at 300 K for InGaN QWs!

The photon flux is

$$\Phi = \eta_{\text{inj}} \eta_i J/q$$

Injection efficiency \equiv capture of carriers by the active region (QWs)

with J the electron current density

LED efficiency

External quantum efficiency (EQE, η): [emitted photons]/[electrons]

$$\eta = \eta_{\text{inj}} \eta_i \eta_{\text{ext}} \longrightarrow \text{extraction efficiency}$$

